

and *not* the hydrogen-bonded *endo*-2'-hydroxyaverrufin (**2b**) as previously assumed.

If a nidurufin-type intermediate is indeed involved in the averrufin to versicolorin A transformation, it seems likely that this intermediate is the *exo*-alcohol nidurufin (**2a**) and not the unknown C-2' *endo* isomer epinidurufin (**2b**), since only **2a** has the favorable stereochemistry for an AR<sub>1</sub>3-promoted rearrangement.

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**Registry No.** **2a**, 28458-23-3; **5**, 518-83-2; **6**, 92762-95-3; **6-ol**, 92762-96-4; **6** (chloro derivative), 92762-97-5; **7**, 92762-98-6; **8**, 92762-99-7; **9**, 92763-00-3; **9a**, 92763-01-4; **10**, 92763-02-5; **11**, 92842-94-9; **5** (methoxymethyl ether), 64517-18-6; **5** (allyl ether), 92763-03-6.

## Spin Multiplet Enhancement in Two-Dimensional Correlated NMR Spectroscopy

Ad Bax\*

Laboratory of Chemical Physics  
National Institute of Arthritis, Diabetes  
and Digestive and Kidney Diseases  
National Institutes of Health  
Bethesda, Maryland 20205

R. Andrew Byrd

Office of Biologics Research and Review  
Center for Drugs and Biologics  
Bethesda, Maryland 20205

A. Aszalos

Food and Drug Administration  
Center for Drugs and Biologics  
Division of Drug Biology, Washington, DC 22064  
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Two-dimensional (2D) homonuclear correlated spectroscopy<sup>1-4</sup> is probably the most widely used 2D NMR experiment to date. It has proven to be a convenient and powerful method for tracing the pattern of homonuclear couplings in molecules with a molecular weight of up to 10000 daltons. However, it appears in practice that the intensity of a cross multiplet, indicating spin coupling, is not directly related to the magnitude of the homonuclear coupling involved. Sometimes protons that have a large scalar interaction show vanishing intensity for the cross multiplet. We propose a method for enhancing selectively the intensity of certain cross peaks of interest and for improving the sensitivity of the experiment.

From the recently introduced operator formalism approach,<sup>5-7</sup> it is easily found that the amount of magnetization transfer,  $R_1$ , from nucleus A to nucleus X in the COSY experiment has the proportionality

$$R_1(t_1) = \sin(\pi J_{AX}t_1) \prod_k \cos(\pi J_{Ak}t_1) \quad (1a)$$

where  $k$  denotes the nuclei other than X to which A is coupled, and  $t_1$  is the duration of the evolution period. Just after the mixing pulse, magnetization transferred from A to X is in antiphase with respect to spin A<sup>2,5</sup> and in phase with respect to other spins,  $n$ , to which X is coupled. The detected X spin magnetization has the  $t_2$  (detection time) dependence:

$$R_2(t_2) = \sin(\pi J_{AX}t_2) \prod_n \cos(\pi J_{Xn}t_2) \quad (1b)$$

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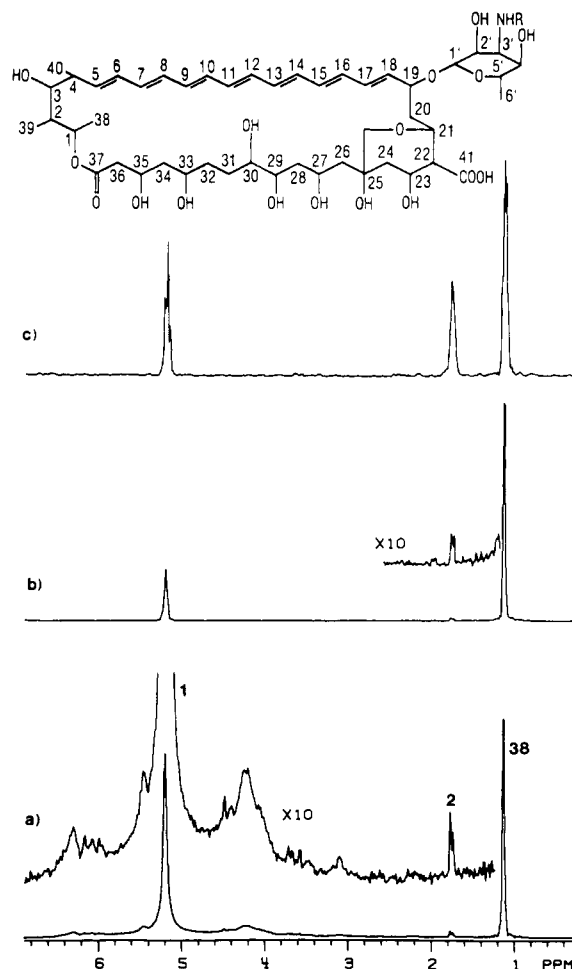
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**Figure 1.** Cross sections parallel to the  $F_1$  axis through the 2D COSY spectra of a sample of amphotericin B, recorded at 500 MHz. The cross sections are taken at the  $F_2$  frequency of proton H1. (a) Cross section if no digital filtering is used, (b) cross section if a single sine bell is used in both dimensions, and (c) the cross section obtained if the optimized filter of eq 3 is used in order to emphasize the H1-H2 cross peak.

If the experiment is performed with phase modulation and  $n$ -type (coherence-transfer echo) selection,<sup>8</sup> the detected X spin magnetization, originating from A, is then given by

$$s_{AX}(t_1, t_2) = M_{0A} [R_1(t_1) R_2(t_2) \exp(-t_1/T_{2A}) \exp(-t_2/T_{2X})] \times \exp(-i\Omega_A t_1) \exp(i\Omega_X t_2) \quad (2)$$

where  $M_{0A}$  denotes the longitudinal A spin magnetization just before the first pulse of the COSY experiment, and  $T_{2A}$  and  $T_{2X}$  are the transverse relaxation times of A and X. From the terms within the square brackets in eq 2, it is seen that the magnetization transfer from A to X does not only depend on the magnitude of  $J_{AX}$  but also on the couplings between A and spins  $k$  and between X and spins  $n$ . Optimal signal to noise for the AX cross multiplet will be obtained for matched filtering, i.e., by multiplying the time domain signal with the function within the square brackets in eq 2. This function is specific for magnetization transferred from A to X and will be nonmatched filtering for other magnetization components. Consequently, diagonal components will be strongly attenuated by such a function, as are other cross peaks that show a different multiplet structure. It is clear from eq 2 that cross peaks between A and X will be hard to observe, if both A and X are coupled to a large number of other spins. However, by the use of a filtering function matched for that particular transfer, all other magnetization components are attenuated relative to the magnetization that contributes to the AX cross peak.

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