## Mixed-time parallel evolution in multiple quantum NMR experiments: Sensitivity

## and resolution enhancement in heteronuclear NMR

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## Supplementary Material

Figure 1 of the main text shows what the 3D MT-PARE-HMQC-NOESY pulse scheme actually "looks like" at different combinations of  $t_1$  and  $t_2$  values, and is conceptually simple. An equivalent scheme, that can be easier to code is shown in Supplementary Material Figure 1S. Actual pulse sequence code for both versions of the pulse program, written for a Bruker Avance system, can be downloaded from http://spin.niddk.nih.gov/bax/pp/.

In the scheme of Figure 1S, the delay  $t_2$  for <sup>15</sup>N chemical shift evolution is always incremented in the outer loop of the pulse code, while the delay  $\Delta_{\rm H}$  is iteratively calculated at each  $t_2$  point during the acquisition according to  $\Delta_{\rm H} = t_2/2$ . In the inner loop of the pulse code (<sup>1</sup>H dimension), the pair of  $180^{\circ} {}^{1}{\rm H}/{}^{15}{\rm N}$  composite pulses is first moved toward time point *a* by incrementing the delay  $t_1{}^{\rm c1}$  (equivalent to  $t_1$  in Figure 1A, main text). Once the delay  $2\delta$  is fully exploited in the above constant-time manner, further <sup>1</sup>H constant-time evolution takes place by incrementing  $t_1{}^{\rm c2}$  (equivalent to  $t_1 - 2\delta$  in Figure 1B, main text), that is, by moving the composite  $180^{\circ} {}^{1}{\rm H}$  pulse during the PARE period toward time point *a*, until  $\Delta_{\rm H}$  cannot be further decremented. Finally, real-time  $t_1$ evolution during  $t_1^{\rm r}$  (equivalent to  $t_1 - 2\delta$  in Figure 1C, main text) follows, if necessary, to fully record <sup>1</sup>H chemical shifts. An advantage of this approach in coding the sequence is that one delay ( $\Delta_{\rm H}$ ) is kept in the inner loop, completely dependent on the other delay ( $t_2$ ) in the outer loop. Note that, although  $t_2$  is incremented linearly, constant-time evolution in  $t_2$  still occurs due to the way this delay is also used for  $t_1$  evolution, in parallel (see Figure 1C, main text and Supplementary Material Figure 2S).

Representative 1D slices of the simulated PARE data taken at the  $t_1$  (or  $t_2$ ) points indicated by colored arrows are shown in B) for the <sup>15</sup>N dimension and in C) for the <sup>1</sup>H dimension in corresponding colors. Note that a weak "ripple", *i.e.* a periodic intensity fluctuation, is seen during the CT evolution (i.e. when  $t_2 < t_1 - 2\delta$ ) in the  $t_2$  (<sup>15</sup>N) dimension, which is absent in the MT-PARE data when recored using the scheme of Figure 1 in the main text (see Figure 3, main text). The ripple results from the fact that incrementation in the  $t_2$  dimension (dwell time 514 µs in our experiment and simulation) can be utilized for constant-time  $t_1$  (dwell time 312 µs) evolution, leaving a remainder of residual time during which relaxation but no chemical shift evolution occurs, and whose exact duration depends on the ratio of the two dwell times. The ripple effect disappears when the dwell times for the two PARE dimensions are chosen to be identical in the two dimensions. In practice, however, the intensity fluctuations fall far below the signal-tonoise and have vanishing effect on the final spectrum, and no limitations need be placed on dwell time selection when the scheme of Figure 1S is used.



Supplementary Material Figure 1S. Alternative version of the 3D MT-PARE-HMQC-NOESY pulse schemes presented in Figure 1 of the main text. Narrow and wide filled bars correspond to 90° and 180° pulses, respectively. Vertically hatched open bars represent composite pulses  $(90^{\circ}_{x}200^{\circ}_{y}90^{\circ}_{x} \text{ for } {}^{1}\text{H} \text{ and } 90^{\circ}_{x}220^{\circ}_{y}90^{\circ}_{x} \text{ for } {}^{15}\text{N})$  employed to reduce RF inhomogeneity and off-resonance effects. Unless otherwise indicated, all pulses have phase x. The constant-time and real-time evolution in MT-PARE is differentiated by the superscripts c and r, respectively, of the incremented delays (i.e. t<sub>1</sub><sup>r</sup>,  $t_1^{c1}$ , and  $t_1^{c2}$  etc.). The 90°  ${}^1H^N$  shaped pulse (labeled as EB) with the EBURP-2 profile,(Geen and Freeman, 1991) has a duration of 1.2 ms (at 800 MHz <sup>1</sup>H frequency), and the pulse labeled as RB is 1.2-ms 180° <sup>1</sup>H<sup>N</sup> shaped pulse with a REBURP profile. These pulses are centered at 8.28 ppm. Broadband <sup>13</sup>C decoupling was achieved using a sequence of adiabatic WURST pulses (Kupče and Freeman, 1995) with a sweep width of 28 kHz and <sup>13</sup>C WURST pulse durations of 10 ms (at 201 MHz <sup>13</sup>C frequency), centered at 117.5 ppm. Phase cycling:  $\phi_1 = -x$ ;  $\phi_2 = x$ , -x; receiver = x, -x. Regular States-TPPI phase cycling of  $\phi_1$  and  $\phi_2$  was used to obtain quadrature detection in the <sup>1</sup>H ( $F_2$ ) and <sup>15</sup>N (*F*<sub>1</sub>) indirect dimensions, respectively. Delay durations:  $\tau = 50$  ms;  $\delta = 4.1$  ms;  $\tau_m = 80$ ms;  $\Delta_{\rm H} = t_2/2$ , i.e. iteratively calculated at each  $t_2$  point during the acquisition. Pulsed field gradients are sine-bell shaped. The gradient pulses have durations of 2 ms and 1 ms, with peak amplitudes of 20 G/cm and 27 G/cm for G<sub>1</sub> and G<sub>2</sub>, respectively.



Supplementary Material Figure 2S. Intensity plot of the 2D time domain data simulated using the alternative MT-PARE pulse scheme shown in Supplementary Material Figure 1S. (B,C) Cross sections taken through the time domain data of (A) at the positions marked by the arrows in (A).